Physics 224 The Interstellar Medium

Lecture #10: Dust Optical Properties, Heating & Cooling

© Karin Sandstrom, UC San Diego - Do not distribute without permission

Outline

- Part I: Dust Optical Properties
- Part II: Dust Heating & Cooling
- Part III: Dust Composition

How we learn about dust

- Extinction: wavelength dependence of how dust attenuates (absorbs & scatters) light
- Polarization: of starlight and dust emission
- Thermal emission from grains
- Microwave emission from spinning small grains
- Depletion of elements from the gas relative to expected abundance
- Presolar grains in meteorites or ISM grains from Stardust mission (7 grains!)

Dust/Light Interaction

Milky Way Dust Extinction Curves





© Karin Sandstrom, UC San Diego - Do not distribute without permission

key reference: Bohren & Huffman textbook



Scattering & absorption result from interaction of grain material with oscillating E & B field

when wavelength of light is < mm magnetic permeability = 1 can ignore magnetic field interaction

plane EM wave $\lambda = 2\pi c/\omega$ E = E₀ e^{ik·r - iωt}

key reference: Bohren & Huffman textbook



plane EM wave $E = E_0 e^{i\mathbf{k}\cdot\mathbf{r} - i\omega t}$ Scattering & absorption result from interaction of grain material with oscillating E & B field

response of material to E field set by *dielectric function*

 $\epsilon(\omega) = \epsilon_1 + i\epsilon_2$

related to refractive index $m=\sqrt{\epsilon}$

Define:

Geometrical Cross Section: πa²

Absorption Cross Section: $C_{abs}(\lambda)$

Scattering Cross Section: $C_{sca}(\lambda)$

Extinction Cross Section: $C_{ext}(\lambda) = C_{abs}(\lambda) + C_{sca}(\lambda)$

Define:

Geometrical Cross Section: πa²

Scattering & Absorption Efficiency Factors

 $Q_{abs} = C_{abs}/\pi a^2$

 $Q_{sca} = C_{sca}/\pi a^2$

Scattering & Absorption by **Small Particles** Scattering Definitions: Albedo = C_{sca}/C_{ext} Differential scattering angle $\frac{dC_{sca}(\theta)}{d\Omega}$

Scattering asymmetry $\langle \cos \theta \rangle = \frac{1}{C_{sca}} \int_0^{\pi} \cos \theta \frac{dC_{sca}(\theta)}{d\Omega} 2\pi \sin \theta d\theta$ factor

- Isotropic scattering $\langle \cos \theta \rangle = 0$
- Forward scattering $\langle \cos \theta \rangle = 1$
- Back scattering $\langle \cos \theta \rangle = -1$



 a/λ - grain size relative to wavelength of light defines different regimes



© Karin Sandstrom, UC San Diego - Do not distribute without permission

$$Q_{\rm abs} = 4 \frac{2\pi a}{\lambda} \operatorname{Im}\left(\frac{\epsilon - 1}{\epsilon + 2}\right) \qquad \qquad Q_{\rm sca} = \frac{8}{3} \left(\frac{2\pi a}{\lambda}\right)^4 \left|\frac{\epsilon - 1}{\epsilon + 2}\right|^2$$

In long wavelength limit, general behavior is:

$$Q_{abs} \sim V/\lambda^2$$
 $Q_{sca} \sim V^2/\lambda^4$

where V = grain volume







© Karin Sandstrom, UC San Diego - Do not distribute without permission

note that at long wavelength: $C_{abs} = Q_{abs}\pi a^2 \propto a^3 \propto m_{dust}$

<u>Absorption</u>

absorption efficiency when $a >> \lambda$ levels off to 1, $C_{abs} = \pi a^2$

for fixed a ←λ <u>Scattering</u> 1 Q_{sca} m=2+iscattering efficiency drops m=1.7+0.1i 0.1 m=1.5+0.01i m=2+0.01i steeply with wavelength m=1.33+0.01i when $a/\lambda << 1$ 0.01 – 0.01

Rayleigh scattering λ^{-4}

 $2\pi |m-1| a/\lambda$

10

10²

0.1

Reflection Nebula vdB1 Image Credit & Copyright: Adam Block, Mt. Lemmon SkyCenter, University of Arizona



© Karin Sandstrom, UC San Diego - Do not distribute without permission



© Karin Sandstrom, UC San Diego - Do not distribute without permission



© Karin Sandstrom, UC San Diego - Do not distribute without permission

Astronomical Dust



© Karin Sandstrom, UC San Diego - Do not distribute without permission

Astronomical Dust



Q_{ext} for astronomical dust analogs



This does not look like the Q_{ext} plots from before - why?

There is a range of grain sizes!

$$\frac{A_{\lambda}}{\text{mag}} = 2.5 \log \left[e^{\tau_{\lambda}}\right] = 1.086 \tau_{\lambda}$$

For a given grain size:

$$\tau_{\nu}(a) = N_d(a) \ Q_{ext}(a) \ \pi a^2$$

Rearrange units to get Weingartner & Draine 2001 eq 7:

$$A(\lambda) = (2.5\pi \log e) \int d\ln a \, \frac{dN_{gr}(a)}{da} \, a^3 Q_{ext}(a, \, \lambda)$$



Continual rise to far-UV means there are more small grains than large grains.





[©] Karin Sandstrom, UC San Diego - Do not distribute without permission

THE SIZE DISTRIBUTION OF INTERSTELLAR GRAINS

JOHN S. MATHIS, WILLIAM RUMPL, AND KENNETH H. NORDSIECK Washburn Observatory, University of Wisconsin-Madison Received 1977 January 24; accepted 1977 April 11

ABSTRACT

The observed interstellar extinction over the wavelength range $0.11 \,\mu\text{m} < \lambda < 1 \,\mu\text{m}$ was fitted with a very general particle size distribution of uncoated graphite, enstatite, olivine, silicon carbide, iron, and magnetite. Combinations of these materials, up to three at a time, were considered. The cosmic abundances of the various constituents were taken into account as constraints on the possible distributions of particle sizes.

Excellent fits to the interstellar extinction, including the narrowness of the $\lambda 2160$ feature, proved possible. Graphite was a necessary component of any good mixture, but it could be used with any of the other materials. The particle size distributions are roughly power law in nature, with an exponent of about -3.3 to -3.6. The size range for graphite is about 0.005 μ m to about 1 μ m. The size distribution for the other materials is also approximately power law in nature, with the same exponent, but there is a narrower range of sizes: about 0.025–0.25 μ m, depending on the material. The number of large particles is not well determined, because they are gray. Similarly, the number of small particles is not well determined because they are in the Rayleigh limit. This power-law distribution is drastically different from an Oort-van de Hulst distribution, which is much more slowly varying for small particles but drops much faster for particles larger than average.

$$Mass(a) \propto \int a^3 \frac{dn}{da} da \propto a^{0.5}$$

Area(a) $\propto \int a^2 \frac{dn}{da} da \propto a^{-0.5}$

most mass in large grains

most area in small grains

 $\frac{dn}{da} \propto a^{-3.5}$

© Karin Sandstrom, UC San Diego - Do not distribute without permission



rate a dust grain of size a absorbs energy

$$\int \left(\frac{dE}{dt}\right)_{abs} = \int h\nu \frac{u_{\nu}}{h\nu} c \ Q_{abs}(\nu)\pi a^2 d\nu$$

$$\int n_{ph} v \ \sigma$$
energy per
absorbed photon

rate a dust grain of size a emits energy

$$\left(\frac{dE}{dt}\right)_{\rm emit} = \int 4\pi B_{\nu}(T) \ Q_{\rm em}(\nu) \ \pi a^2 d\nu$$

blackbody emitting over 4π str with efficiency Q_{em}

rate a dust grain of size a absorbs energy

$$\left(\frac{dE}{dt}\right)_{\rm abs} = \int h\nu \ \frac{u_{\nu}}{h\nu} \ c \ Q_{\rm abs}(\nu)\pi a^2 d\nu$$

$$\underline{\text{LTE}} \qquad u_{\nu} = \frac{4\pi}{c} B_{\nu}(T)$$
$$\left(\frac{dE}{dt}\right)_{\text{abs}} = \left(\frac{dE}{dt}\right)_{\text{emit}}$$

© Karin Sandstrom, UC San Diego - Do not distribute without permission

in

rate a dust grain of size a absorbs energy

$$\left(\frac{dE}{dt}\right)_{\rm abs} = \int h\nu \ \frac{u_{\nu}}{h\nu} \ c \ Q_{\rm abs}(\nu)\pi a^2 d\nu$$

<u>in LTE</u>

$$\int \frac{4\pi}{c} B_{\nu}(T) \ c \ Q_{\rm abs}(\nu) \pi a^2 d\nu = \int 4\pi B_{\nu}(T) \ Q_{\rm em}(\nu) \pi a^2 d\nu$$

Therefore:
$$Q_{abs} = Q_{em}$$

rate a dust grain of size a absorbs energy

$$\left(\frac{dE}{dt}\right)_{\rm abs} = \int h\nu \ \frac{u_{\nu}}{h\nu} \ c \ Q_{\rm abs}(\nu)\pi a^2 d\nu$$

Define "spectrum averaged absorption cross section"

$$\left\langle Q_{\rm abs}\right\rangle_* \equiv \frac{\int u_{*\nu} Q_{\rm abs}(\nu) d\nu}{\int u_{*\nu} d\nu}$$

so that:
$$\left(\frac{dE}{dt}\right)_{abs} = \langle Q_{abs} \rangle_* \pi a^2 u_* c$$



 $< Q_{abs} > * for$ the average interstellar radiation field in the MW, and two astronomical dust analogs.

rate a dust grain of size a emits energy

$$\left(\frac{dE}{dt}\right)_{\text{emit}} = \int 4\pi B_{\nu}(T) \ Q_{\text{em}}(\nu) \ \pi a^2 d\nu$$
$$Q_{\text{abs}} = Q_{\text{em}}$$

Define "Planck averaged emission efficiency"

$$\langle Q_{\rm abs} \rangle_{\rm T} \equiv \frac{\int B_{\nu}(T) Q_{\rm abs} d\nu}{\int B_{\nu}(T) d\nu}$$

so that:
$$\left(\frac{dE}{dt}\right)_{\rm em} = 4\pi a^2 \langle Q_{\rm abs} \rangle_{\rm T} \sigma T^4$$



[©] Karin Sandstrom, UC San Diego - Do not distribute without permission

In steady state, emission = absorption.

$$\left(\frac{dE}{dt}\right)_{\rm abs} = \langle Q_{\rm abs} \rangle_* \pi a^2 \ u_* \ c$$
$$\left(\frac{dE}{dt}\right)_{\rm em} = 4\pi a^2 \ \langle Q_{\rm abs} \rangle_{\rm T} \ \sigma T^4$$

for MW interstellar radiation field and dust properties we found: $\langle Q_{\rm abs} \rangle_* \sim 0.8 (a/0.1 \mu m)^{0.85}$ carbon $\langle Q_{\rm abs} \rangle_* \sim 0.18 (a/0.1 \mu m)^{0.6}$ silicate

© Karin Sandstrom, UC San Diego - Do not distribute without permission

In steady state, emission = absorption.

$$\left(\frac{dE}{dt}\right)_{\rm abs} = \langle Q_{\rm abs} \rangle_* \pi a^2 \ u_* \ c$$
$$\left(\frac{dE}{dt}\right)_{\rm em} = 4\pi a^2 \ \langle Q_{\rm abs} \rangle_{\rm T} \ \sigma T^4$$

for MW interstellar radiation field and dust properties we found: $\langle Q_{\rm abs} \rangle_T \sim 8 \times 10^{-6} T^2$ carbon $\langle Q_{\rm abs} \rangle_T \sim 1.3 \times 10^{-5} T^2$ silicate

© Karin Sandstrom, UC San Diego - Do not distribute without permission

Very weak dependence of equilibrium temperature on grain size.

$$T \approx 22.3 (a/0.1 \mu m)^{-1/40} U^{1/6} K$$

silicate

carbon

 $T \approx 16.4 (a/0.1 \mu m)^{1/15} U^{1/6} K$