#### Physics 224 The Interstellar Medium

Lecture #6: Radiative Transfer

- Part I: Basics of Radiative Transfer
- Part II: HI 21-cm Radiative Transfer
- Part III: Absorption Lines

Motions of individual particles

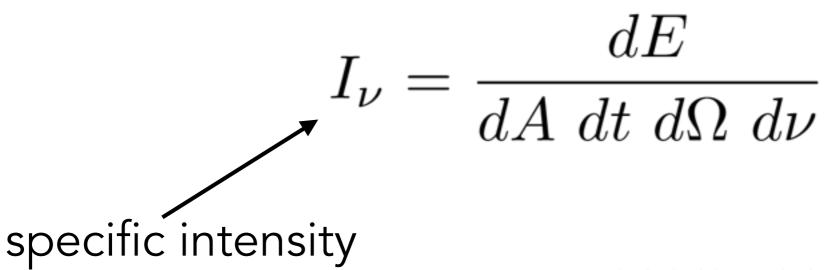
On scales > mean free path for collisions Fluid dynamics Propagation of individual photons

On scales ≫ λ ↓ Radiative Transfer

Transport Phenomena: <u>https://en.wikipedia.org/wiki/Transport\_phenomena</u>

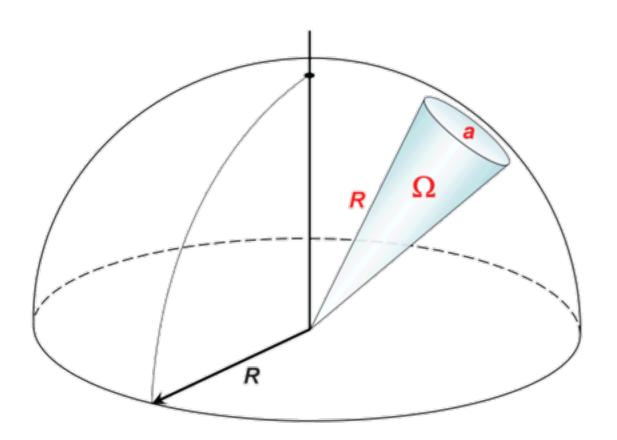
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Description of the radiation field in full detail:



energy per unit time (dt) per unit frequency (dv) passing through area (dA) from solid angle (dΩ)

Solid Angle

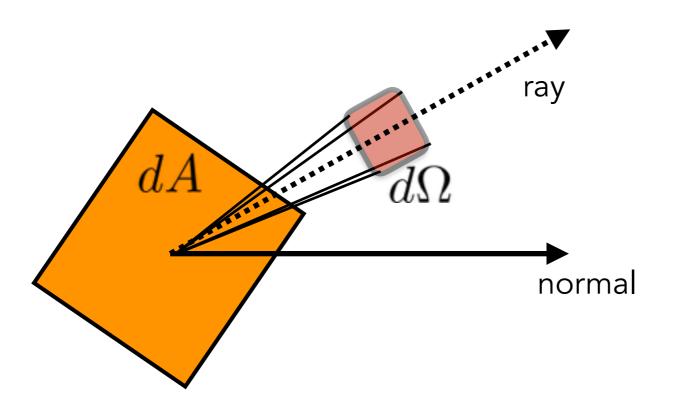


Two-dimensional angle subtended by an object.

Unit = steradian (sr)

 $4\pi$  sr in a sphere

Solid angle specifies direction

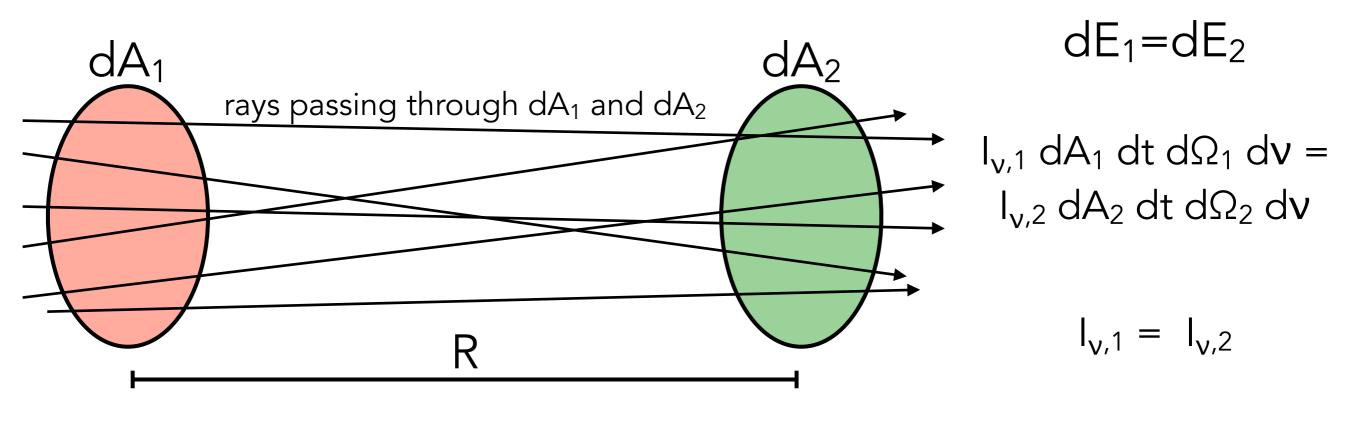


Units of Specific Intensity

 $[I_v] = \text{energy} (\text{time})^{-1} (\text{area})^{-1} (\text{solid angle})^{-1} (\text{frequency})^{-1}$ 

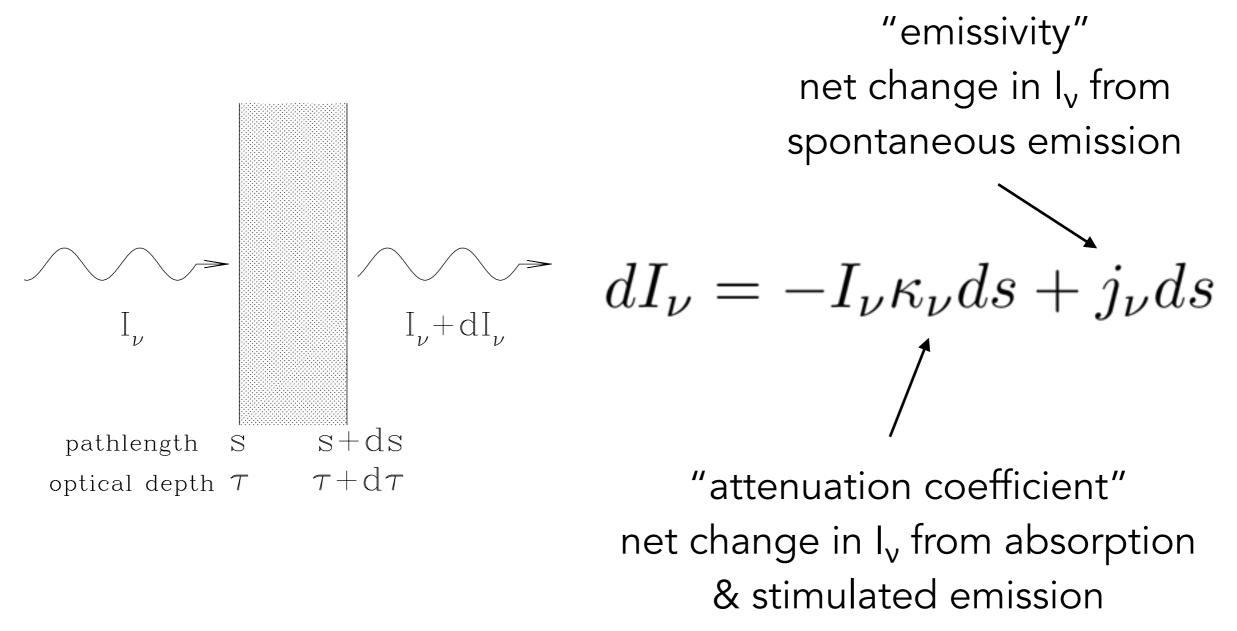
$$[I_v] = ergs s^{-1} cm^{-2} sr^{-1} Hz^{-1}$$

When there is no emission/absorption, intensity along a ray is constant.



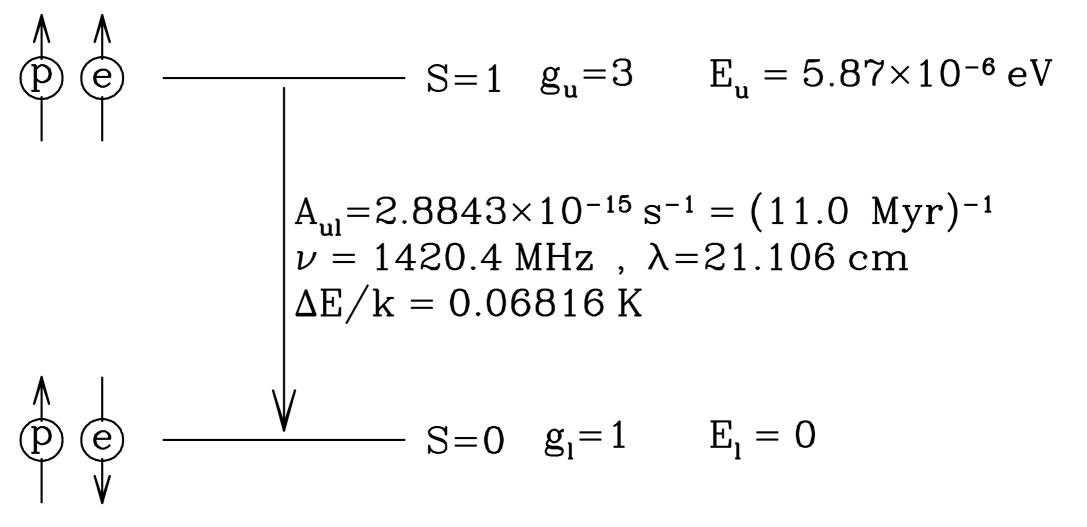
 $d\Omega_1 = \text{solid}$  angle subtended by  $dA_2$  at  $dA_1 = dA_2/R^2$   $I_v = \text{constant}$  $d\Omega_2 = \text{solid}$  angle subtended by  $dA_1$  at  $dA_2 = dA_1/R^2$ 

Equation of Radiative Transfer ignoring scattering



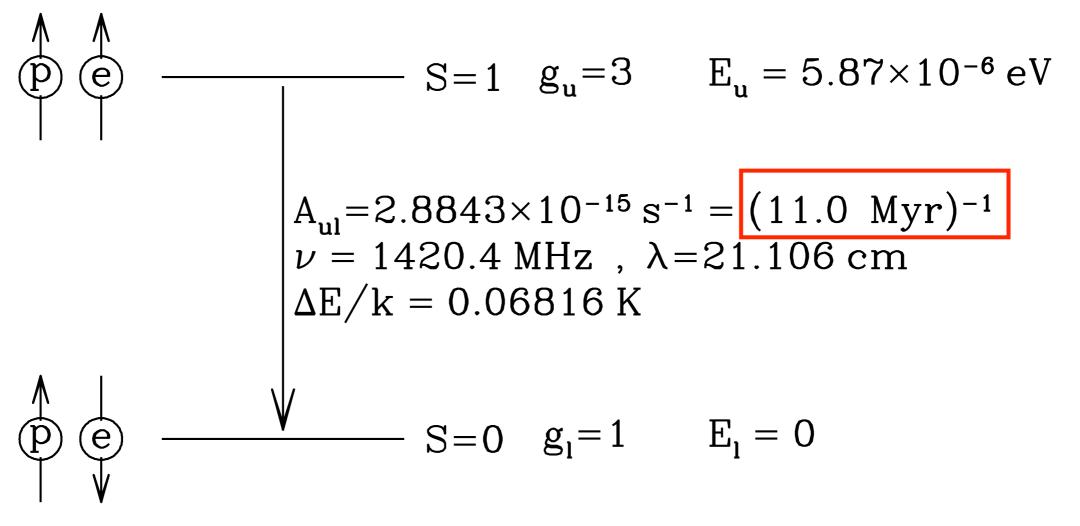
# Part II: HI 21-cm Radiative Transfer

Hyperfine splitting due to interaction of electron spin and nuclear spin.



\*remember: degeneracy = 2(S+1) = 3 for upper

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 $T_{exc} \equiv T_{spin} \gg 0.0682 \text{ K}$ 

because cosmic microwave background can populate levels

$$\frac{n_u}{n_l} \equiv \frac{g_u}{g_l} e^{-h\nu_{ul}/kT_{exc}} = 3e^{-0.0682/T_{spin}} \approx 3$$

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Under all ISM conditions, 75% of HI is in upper level. Emissivity is independent of  $T_{spin}$ !!

$$j_{\nu} = n_u \frac{A_{ul}}{4\pi} h \nu_{ul} \phi_{\nu} = \frac{3}{16\pi} A_{ul} h \nu_{ul} n(\text{H I}) \phi_{\nu}$$

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$$\kappa_{\nu} = \frac{h\nu}{4\pi} n_l B_{lu} \phi_{\nu} \left( 1 - e^{-E_{ul}/kT_{exc}} \right)$$

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from last week:  

$$\frac{g_l B_{lu}}{g_u B_{ul}} = 1$$

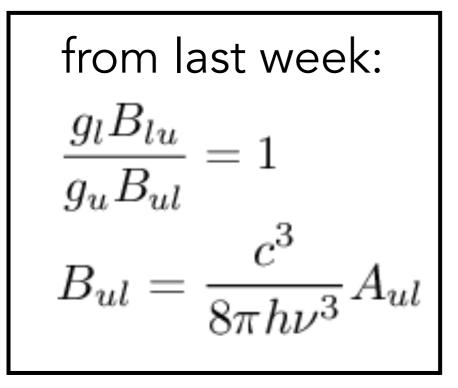
$$B_{ul} = \frac{c^3}{8\pi h\nu^3} A_{ul}$$

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use:  $E_{ul} \ll kT_{spin}$ 



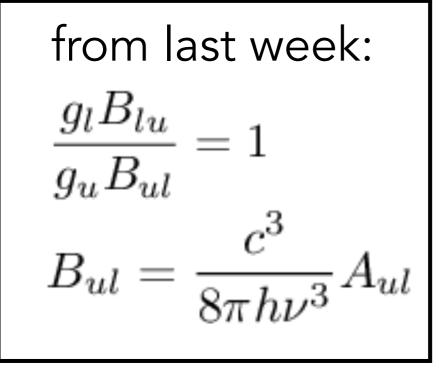
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to get: 
$$e^{-E_{ul}/kT_{spin}} = 1 - h\nu_{ul}/kT_{spin}$$



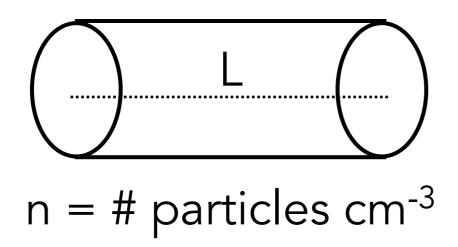
put all this together and we find:

$$\kappa_{\nu} \approx \frac{3}{32\pi} A_{ul} \frac{hc\lambda_{ul}}{kT_{spin}} n(\text{H I})\phi_{\nu}$$

absorption coefficient depends inversely on T<sub>spin</sub> this is a consequence of <u>stimulated emission</u> <u>not being negligible!</u>

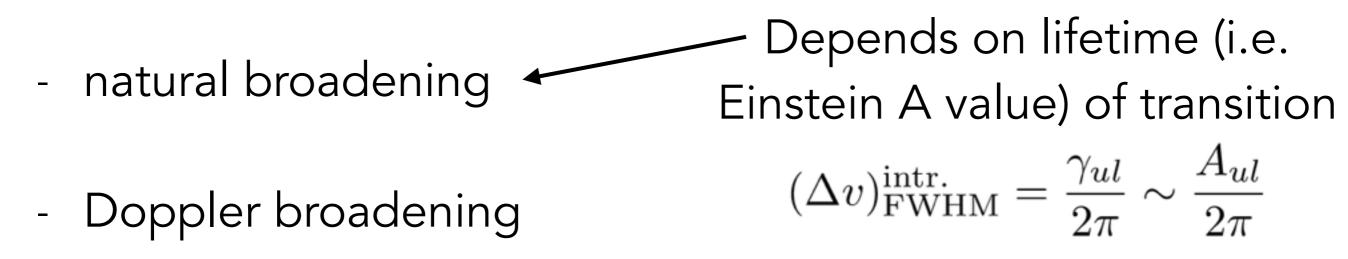
# Column Density $N(H I) \equiv \int ds \ n(H I)$

for uniform volume density N = nL



column density is number of particles per unit area along a path length L

recall from last week, two parts of line profile  $\phi_{v}$ :



Because lifetime of 21-cm transition is SO LONG natural broadening is negligible.

Line profile depends on velocity dispersion.

$$\tau_{\nu} = 2.19 \left( \frac{N(\text{H I})}{10^{21} \text{cm}^{-2}} \right) \left( \frac{T_{spin}}{100K} \right)^{-1} \left( \frac{\sigma_{v}}{\text{kms}^{-1}} \right)^{-1} e^{-u^{2}/2\sigma_{v}}$$

u = velocity difference from line center

Lots of regions where optical depth in HI can be significant.

In the optically thin case  $\tau_v \ll 1$ 

$$I_{\nu} = I_{\nu}(0) + \int j_{\nu} ds = I_{\nu}(0) + \frac{3}{16\pi} A_{ul} h\nu_{ul} \phi_{\nu} N(\text{H I})$$

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Say 
$$I_v(0) = 0$$

$$\int I_{\nu} d\nu = \frac{3}{16\pi} A_{ul} h\nu_{ul} N(\text{H I})$$
  
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Note: radio astronomers often convert  $I_v$  to brightness temperature